Existence of Chern-Simons modified gravity based on discovery of many galaxies without dark matter

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ABSTRACT

The existence of dark matter has been assumed to explain the flat rotation curves of galaxies. Recently, however, many studies have reported galaxies without dark matter. Therefore, we need to reconsider dark matter theory. Although modified Newtonian dynamics (MOND) (which is another candidate to explain flat rotation curves without dark matter) can describe spiral galaxies well, it is inapplicable to recent observations. Here, we focus on considering the rotational motion of galaxies on the basis of Newtonian dynamics with a relativistic correction, and we find a new empirical formula for flat rotation curves. The simple formula is derived from Chern-Simons (CS) modified gravity. We may be able to understand the dynamics of all galaxies since the helicity C_2 in CS theory can take different values. In this study, we calculate CS gravitational acceleration and compare it with MOND gravitational acceleration. Furthermore, we compute the log variance ratio between CS data and MOND data. Consequently, we discover that CS data fits the observed gravitational acceleration better than MOND data for a large radius where flat rotation curves appear. It also means that we can explain the flat rotation curves of galaxies without dark matter. Rotational dynamics regarded as MOND or dark matter phenomena may be interpreted as CS vortex phenomena.

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Key words: Chern-Simons modified gravity - modified Newtonian dynamics - dark matter

1 1 INTRODUCTION

There is a mystery regarding the flat rotation curves of galax-²⁸ 2 ies. Galaxies rotate as if they contain an invisible mass. To 29 3 explain this phenomenon, the existence of dark matter has ³⁰ 4 been assumed (Rubin et al. 1978). Although X-rays observed ³¹ in an elliptical galaxy offer its evidence, a black hole may ³² 6 be an X-ray source in such a galaxy (Haaften et al. 2019). ³³ 7 In terms of gravitational lensing in a cluster of galaxies, we ³⁴ 8 cannot decide whether it is a proof of dark matter because 35 g the amount of dark matter estimated by gravitational lensing ³⁶ 10 was frequently less than expected (Meneghetti et al. 2020). In ³⁷ 11 addition, many galaxies without dark matter have been dis- ³⁸ 12 covered (Dokkum et al. 2018; Gui et al. 2020; Montes et al. 39 13 2020). Thus, we need to reconsider dark matter theory (Lelli ⁴⁰ 14 41 15 2014; Firmani et al. 2001). Modified Newtonian dynamics (MOND) (Milgrom 1983; Mc- ⁴² 16 Gaugh 2008; Smolin 2017) may provide another way to solve 43 17 this fundamental problem regarding the universe. Although 44 18 MOND is a phenomenological theory without dark matter 45 19 (McGaugh & Lelli 2016; McGaugh 2008), it is inapplicable to ⁴⁶ 20 non-rotating elliptical galaxies. In contrast, we are interested 47 21 in Chern-Simons (CS) modified gravitational theory (Jackiw 22 & Pi 2003; Konno et al. 2007, 2009; Alexander & Yunes 2009), 23 which deals with a vortex system including a singularity. In 24 this study, we discovered a new empirical formula which de-25

scribes flat rotation curves. The formula corresponds with the helicity C_2 of the space-time structure in CS theory. Hereafter, we use CS^{*} to indicate the empirical CS. In principle, CS^{*} may be able to describe flat rotation curves for any type of galaxy from elliptical to spiral since C_2 can take from zero to an arbitrary value (Konno et al. 2008). We show that CS^{*} is the most suitable theory for understanding galactic dynamics.

By using the new empirical formula, we calculated gravitational acceleration and fitted observed data for spiral galaxies. Moreover, we computed the log variance ratio between CS* and MOND, $k = \delta_{\rm CS}^2 / \delta_{\rm MOND}^2$, for small values of gravitational acceleration from baryonic matter: $g_{\rm bar} \leq 10^{-10} {\rm m/s}^2$ (that is, $r > 5 \times 10^{21}$ m in the case of the Milky Way galaxy). We show that CS* can be applied to the problem of flat rotation curves. This is consistent with the discovery of galaxies without dark matter.

This paper is organized as follows. In §2, we outline the limitations of dark matter and MOND. In §3, we describe the relationship between the new empirical formula and CS theory. Each theory is compared in §4, which comprises the discussion and conclusion.

48 2 DARK MATTER AND MOND

Dark matter is one of the strong candidates for explaining 49 flat rotation curves. In dark matter theory, spiral galaxies 50 contain several times more dark matter than luminous (bary-51 onic) matter. However, it is currently unclear whether or not 52 this theory is applicable to elliptical galaxies (Romanowsky 53 et al. 2003). For example, the deprivation of dark matter was 54 recently observed in the elliptical galaxy NGC 7507 (Lane 55 56 et al. 2015). In addition, many elliptical galaxies without dark matter have been discovered (Dokkum et al. 2018; Gui et al. 57 2020). Further studies may solve this problem, the depriva-58 tion or the absence of dark matter (Bidin et al. 2012, 2015). 59 In contrast, MOND is a phenomenological theory that can 60 also describe the flat rotation curves of spiral galaxies. This 61 model was first proposed by Milgrom (Milgrom 1983), and 62 McGaugh developed the following gravitational acceleration 63 formula (McGaugh & Lelli 2016; McGaugh 2008) : 64

$$g_{\text{MOND}} = \frac{g_{\text{bar}}}{1 - e^{-\sqrt{g_{\text{bar}}/g_{\dagger}}}} \tag{1}$$

where g_{\dagger} is a universal acceleration constant and g_{bar} is 65 gravitational acceleration from baryonic matter. Figure 1-(a) 66 shows the results obtained with this equation and the ob-67 served gravitational acceleration of spiral galaxies. We theo-68 retically consider MOND data (McGaugh & Lelli 2016). In 69 Figure 1-(a), MOND provides a good description of flat ro-70 tation curves. Equation (1) must also be applicable to (non-71 rotating) elliptical galaxies because it is a universal modifi-72 cation of Newtonian dynamics. This universality contradicts 73 recent discoveries of galaxies without dark matter. 74 Therefore, although dark matter theory and MOND are gen-75 erally used to explain flat rotation curves, they may not pro-76 vide information about elliptical galaxies (except for certain 77

rotating galaxies). Here, we describe flat rotation curves from
Newtonian dynamics with a relativistic correction. In the
next section, we show that a new empirical formula is applicable to various galaxies because of a relationship with CS
theory.

83 3 NEW EMPIRICAL FORMULA

In this study, we newly added a function to Newton's equation of motion. When we consider rotational galactic motion
based on this expansion, the next equation of rotational speed
v_{rot} is assumed

$$v_{\rm rot} = \sqrt{\frac{GM}{r}} + \frac{C_{\rm rot}}{2} \tag{2}$$

where M is the mass of a galaxy, G is the universal gravitational constant and r is the radius from the center of a galaxy. The first term indicates rotational speed derived from Newtonian mechanics. We assumed that $C_{\rm rot}$ is a function of r. We analyzed flat rotation curves with this equation. First, we assumed a standard Laurent expansion for $C_{\rm rot}$:

$$C_{\rm rot} = \sum_{n=-\infty}^{\infty} a_n r^n \tag{3}$$



Figure 1. (a) The horizontal axis represents g_{bar} in log scale and the vertical axis represents g_{obs} or g_{MOND} in log scale. The red line shows g_{obs} , which is the observed gravitational acceleration of spiral galaxies. The blue line shows g_{MOND} , which was calculated with Equation (1). (b) The horizontal axis represents g_{bar} in log scale and the vertical axis represents $g_{\rm obs}$ or $g_{\rm CS^*}$ in log scale. The red line shows g_{obs} as in Figure 1-(a). The blue line shows g_{CS^*} , which we calculated with Equation (4), (5) and (6). Each figure shows that MOND and CS* can describe flat rotation curves since each set of calculated data is comparable to observed data. In particular, we focused on seven points from a long distance and confirmed the difference between MOND and CS* in more detail. In the insets, the horizontal axis represents g_{bar} in log scale and the vertical axis represents $g_{\text{MOND}} - g_{\text{obs}}$ or $g_{\text{CS}^*} - g_{\text{obs}}$, where $g_{\rm MOND}-g_{\rm obs}$ and $g_{\rm CS}{}^*-g_{\rm obs}$ are the difference from observed data represented by a thick black line. By computing the log variance ratio between CS^{*} and MOND, $k = \delta_{CS^*}^2 / \delta_{MOND}^2$, based on the result, we obtained k = 0.885. This indicates that CS^{*} fits the observed data better than MOND for a large radius.

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As explained below, we discovered that the following three138

terms are effective for flat rotation curves:

To verify the validity of our fitting, we calculated the log variance defined by:

$$C_{\rm rot} = \alpha r + \frac{\beta}{r} + \gamma \tag{4}$$

where α , β and γ are constant and expressed by a_n in Equation (3). We found $\alpha = 2.4 \times 10^{-17} \text{ s}^{-1}$, $\beta = -2.1 \times 10^{26}$

 m^2/s , $\gamma = 3.5 \times 10^5 m/s$ for the best fitting. γ is the most₁₄₀ 98 effective term for $C_{\rm rot}$ because r and r^{-1} work as correction₁₄₁ 99 terms: these terms correct at a larger radius and a smaller 142 1 00 radius, respectively. In fact, the contribution of the r^{-1} term₁₄₃ 1 01 is 1 to 2% of γ , thus the new empirical formula is just the₁₄₄ 102 Laurent expansion around γ up to the first order in r. If we₁₄₅ 103 compare space-time and fluid, we can treat αr as a forced₁₄₆ 104 vortex, and βr^{-1} as a free vortex. These two terms are inter-1 05 preted as the conservation of angular momentum. γ is linked 148 106 to a topological term. The most important γ term in Equa-149 107 tion (4) is derived from CS theory for the following reasons. $_{150}$ 1 08 CS modified gravity in (3 + 1) dimensions was proposed by 151 109 Jackiw and Pi (Jackiw & Pi 2003). The effects of the CS₁₅₂ 110 action in gravitational fields have been discussed in several₁₅₃ 111 studies (Alexander & Yunes 2007a,b). A perturbative study₁₅₄ 112 with small angular momentum provides the following velocity $_{155}$ 113 (Konno et al. 2008): 114 156

$$v_{\rm CS} = \sqrt{\frac{GM}{r}} + \frac{C_2}{2} + \frac{r_s C_2}{r} + \frac{r_s^2}{2r^2} \left[C_1 + C_2 + 2C_2 \log\left(\frac{r - r_s}{r_s}\right)_{\rm rg}^{15} \right]$$
(5) 163

162 where $r_s = \frac{2GM}{c^2}$ is the Schwarzschild radius, M is the mass₁₆₃ 115 of a galaxy, G is the universal gravitational constant, c is the $_{164}$ 116 speed of light, r is the radius of a galaxy and C_1 , C_2 are the₁₆₅ 117 helicity in CS. At a large r, Equation (5) simply becomes₁₆₆ 118 Equation (2) and $C_{\rm rot}$ corresponds to the helicity C_2 , which 167 119 is an important term in CS. This implies that our assumed $_{168}$ 120 Equation (2) is strongly related to CS theory. 121 169 Gravitational acceleration of CS^* is calculated by the follow-170 122

¹²² Gravitational acceleration of CS is calculated by the follow-₁₇₀ ¹²³ ing equation: 171

$$g_{\rm CS^*} = \frac{v_{\rm CS^*}^2}{(6)_{174}}$$

$$y_{\rm S}^* = \frac{1}{r}$$
 (0)174

Figure 1-(b) shows the results of CS^{*} gravitational accelera-176 124 tion with Equation (6). CS^* can describe flat rotation curves₁₇₇ 125 as well as MOND. In particular, the results were comparable₁₇₈ 126 to observed data when we added the r^{-1} term. This indi-179 127 cates the importance of the r^{-1} correction term to explain₁₈₀ 128 flat rotation curves. CS* can be applied to various galaxies₁₈₁ 129 since the helicity C_2 in Equation (5) can take different values₁₈₂ 1 30 including zero. We can understand the motion of galaxies by 183 1 31 considering CS gravity. 1 32 184

133 4 DISCUSSION AND CONCLUSION

We compared CS* and MOND by computing the differences190
between both theories and observed data. Those results are191
shown in the insets in Figure 1. Gravitational acceleration192
values based on CS* and MOND fit well at a large radius.193

$$\delta_{\rm th}^2 = \frac{1}{N} \sum_{n=1}^{N} \left(\log g_{\rm obs} - \log g_{\rm th} \right)^2 \tag{7}$$

where $g_{\rm th}$ denotes $g_{\rm CS^*}$ or $g_{\rm MOND}$, and N is number of data. Then, we calculated the log variance ratio between CS^{*} and MOND, $k = \delta_{\rm CS^*}^2 / \delta_{\rm MOND}^2$, for small values of $g_{\rm bar} \leq 10^{-10}$ m/s² (that is, $r > 5 \times 10^{21}$ m in the case of the Milky Way galaxy). As a result, we obtained $k = \delta_{\rm CS^*}^2 / \delta_{\rm MOND}^2 = 0.885$. The log variance ratio increases for a smaller radius. However, this deviation is not important here because flat rotation curves appear for galaxies with larger radiuses. That is, the result indicates that a CS^{*} fit can describe flat rotation curves better than a MOND fit. Moreover, CS^{*} has a decisively different feature, namely it can be applied to elliptical galaxies. In that situation, C_2 in CS can take zero or a small value. CS^{*} may be a more general theory that solves the problem with MOND.

Figure 2 outlines the differences between dark matter theory, MOND and CS^{*}. Dark matter theory and MOND give us information about the flat rotation curves of spiral galaxies. However, if they were the only theories that can be used to explain flat rotation curves, then paradoxically the dark matter deficiency in elliptical galaxies favors dark matter theory. This strange logic fails if we further consider CS^{*}. In principle, CS^{*} can describe not only spiral galaxies but also elliptical galaxies since the helicity C_2 can take different values. We may be able to classify galaxies based on Hubble classification with CS^{*}. The r^{-1} term in C_2 that we discovered may be computed from a different form of CS, e.g. dynamical CS (Konno & Takahashi 2014; Yagi et al. 2012). Although it is difficult to find an exact solution, we plan to work on this problem in a future study.

In spite of many efforts, there has as yet been no direct laboratory detection of dark matter particles. On the other hand, if we change our viewpoint and assume that CS helicity exists in the universe, we can obtain a compatible explanation for the flat rotation curves of galaxies. In this study, we confirmed that CS^{*} can describe spiral galaxies as well as MOND without dark matter. With a large radius (where flat rotation curves appear) the log variance ratio k = 0.885 implies that CS^{*} fits observed gravitational acceleration better than MOND. In addition, there is a definite difference between CS* and MOND. CS, which is related to the new empirical formula, is a universal theory and can be applied to all galaxies. This theory may provide us some solutions for flat rotation curves. CS structure has the potential to solve the problems posed by many astrophysical phenomena (Canizares et al. 2012; Barrientos et al. 2019) whereas dark matter or MOND is not required in some cases (Lane et al. 2009). Ultimately, it can be applied to condensed matter physics and particle physics (Cisterna et al. 2019). We plan to use CS to research many problems including exploring axial symmetry theory (Bahamonde et al. 2021).



Figure 2. Classification of dark matter, MOND and CS* based²⁴¹ on Hubble classification. Arrows with solid lines indicate the re-²⁴² gion in which each model is applicable. The dashed lines indicate²⁴³ the region where each theory may not be applicable. We do not²⁴⁴ currently know why galaxies rotate so strangely (winding dilemma²⁴⁵ etc.). MOND can explain flat rotation curves only for spiral galax-²⁴⁶ ies and some rotating elliptical galaxies (Samurović & Vudragović²⁴⁷ 2019), that is, this theory is not applicable to non-rotating galaxies.²⁴⁸ Dark matter theory is a strong candidate for solving the problem of²⁴⁹ flat rotation curves. Recent observations for elliptical galaxies indi-²⁵⁰ cate the existence of galaxies with less dark matter than expected.²⁵¹ On the other hand, we can understand all galaxies with CS* since²⁵² universal CS theory has the important helicity C_2 . Each galaxy²⁵³ may have different (including vanishing) values of C_2 depending on its type. Credit: Ville Koistinen.

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